

# Stopping power of electron gas

Yu.V. Arkhipov<sup>1</sup>, A.B. Ashikbayeva<sup>1</sup>, A. Askaruly<sup>1</sup>, A.E. Davletov<sup>1</sup>, D.Yu. Dubovtsev<sup>1</sup>, S.A. Syzganbayeva<sup>1</sup>, I.M. Tkachenko<sup>2</sup>

<sup>1</sup> Al-Farabi Kazakh National University, Almaty, KAZAKHSTAN

<sup>2</sup> Universitat Politècnica de València, Valencia, SPAIN

The polarizational stopping power of heavy ions in fully ionized plasmas is described by the Lindhard formula [1]:

$$\left[ -\frac{dE}{dx} \right]^{pol} = \frac{2(Z_p e)^2}{\pi v^2} \int_0^\infty \frac{dk}{k} \int_0^{kv} \omega^2 \left( -\text{Im} \left[ \frac{\varepsilon^{-1}(k, \omega)}{\omega} \right] \right) d\omega,$$

whose high-velocity asymptotic form was found by Bohr, Bethe, and Larkin [2]:

$$\left[ -\frac{dE}{dx} \right]_{v \rightarrow \infty}^{pol} \cong \left( \frac{Z_p e \omega_p}{v} \right)^2 \ln \left[ \frac{2m_e v^2}{\hbar \omega_p} \right].$$

The plasma inverse dielectric function (IDF),  $\varepsilon^{-1}(k, \omega)$ , was determined in [3] within the moment approach [4], complemented by some physical observations in terms of two characteristic frequencies  $\omega_1$  and  $\omega_2$ , which are the ratios of the frequency power moments of the IDF imaginary part, in the following form:

$$\varepsilon^{-1}(k, \omega) = 1 + \frac{\omega_p^2 (\sqrt{2}\omega_1 \omega + i\omega_2^2)}{\sqrt{2}\omega_1 \omega (\omega^2 - \omega_2^2) + i\omega_2^2 (\omega^2 - \omega_1^2)}.$$

The frequencies  $\omega_1$  and  $\omega_2$  can be rigorously evaluated using the static structure factor (SSF) of the system. Nevertheless, here we employ the following interpolating expressions [5,6]:

$$\omega_1^2 = \omega_1^2(k) = \omega_p^2 \left( 1 + k^2 k_D^{-2} + k^4 k_q^{-4} \right), \quad \omega_2^2 = \omega_2^2(k) = \omega_p^2 \left( 1 + \frac{\langle v_e^2 \rangle k^2}{\omega_p^2} - \frac{v_{int}^2 k^2}{\omega_p^2} \right),$$

The interpolation and fitting parameters introduced are chosen as follows:

$$v_{int}^2 = -\frac{4}{15} \frac{\Gamma^{3/2}}{\beta m_e} \left( \frac{-0.9052}{\sqrt{0.6322 + \Gamma}} + \frac{0.27243}{1 + \Gamma} \right), \quad \Gamma = \beta e^2 / a, \quad k_q^4 = 12r_s / a^4, \quad r_s = a / a_B,$$

$a$  and  $a_B$  are the Wigner-Seitz and Bohr radii, respectively,  $\beta = 1/(k_B T)$ ,  $k_D^{-1}$  is the Debye radius,  $k_B$  stands for the Boltzmann constant with  $T$  being the plasma temperature.

The numerical results obtained for the energy losses of heavy ions moving in an electron gas are found in good agreement with the PIC simulation data [7].

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